

**Two-fluid dark matter models**Tiberiu Harko<sup>\*</sup>*Department of Physics and Center for Theoretical and Computational Physics, The University of Hong Kong, Pok Fu Lam Road, Hong Kong, People's Republic of China*Francisco S. N. Lobo<sup>†</sup>*Centro de Astronomia e Astrofísica da Universidade de Lisboa, Campo Grande, Ed. C8 1749-016 Lisboa, Portugal*

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We investigate the possibility that dark matter is a mixture of two noninteracting perfect fluids, with different four-velocities and thermodynamic parameters. The two-fluid model can be described as an effective single anisotropic fluid, with distinct radial and tangential pressures. The basic equations describing the equilibrium structure of the two-fluid dark matter model, and of the tangential velocity of test particles in stable circular orbits, are obtained for the case of a spherically symmetric static geometry. By assuming a nonrelativistic kinetic model for the dark matter particles, the density profile and the tangential velocity of the dark matter mixture are obtained by numerically integrating the gravitational field equations. The cosmological implications of the model are also briefly considered, and it is shown that the anisotropic two-fluid model isotropizes in the large-time limit.

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**I. INTRODUCTION**

Cosmological observations provide compelling evidence that about 95% of the content of the Universe resides in two unknown forms, denoted dark matter and dark energy, respectively. The former resides in bound objects as non-luminous matter, and the latter in the form of a zero-point energy that pervades the whole Universe [1]. Dark matter is thought to be composed of cold neutral weakly interacting massive particles, beyond those existing in the standard model of particle physics, and not yet detected in accelerators or in direct and indirect searches. There are many possible candidates for dark matter, the most popular ones being the axions and the weakly interacting massive particles (WIMP) (for a review of the particle physics aspects of dark matter see [2]). Their interaction cross section with normal baryonic matter, while extremely small, are expected to be nonzero and we may expect to detect them directly. Scalar fields or other long range coherent fields coupled to gravity have also been used intensively to model galactic dark matter [3]. Recently, the possibility that the galactic dynamics of massive test particles may be understood in the context of modified theories of gravity, without the need for dark matter, was also explored [4].

Despite these important achievements, at galactic scales  $\sim 10$  kpc, the  $\Lambda$ CDM model meets with severe difficulties in explaining the observed distribution of the invisible matter around the luminous one. In order for galactic gas and stars to attain a constant velocity in the outer regions, the density profile of the dark matter must fall approximately with  $r^{-2}$  in the respective regions. These

profiles resemble the density structure of an isothermal and self-gravitating system of particles, characterized by a constant velocity dispersion. In fact,  $N$ -body simulations, performed in this scenario, predict that bound halos surrounding galaxies must have very characteristic density profiles that feature a well pronounced central cusp [5]. On the observational side, high-resolution rotation curves show, instead, that the actual distribution of dark matter is much shallower than the above, and it presents a nearly constant density core [6].

The possibility that dark matter could be described by a fluid with nonzero effective pressure was also considered in the physical literature [7]. In particular, in [8] it was assumed that the equation of state of the dark matter halos is polytropic. The fit with a polytropic dark halo improves the velocity dispersion profiles.

Because of their importance in astrophysics and cosmology, two-fluid models have been extensively investigated in the physical literature. In particular, dust shells of finite size, moving with different four-velocities and forming a common region, where the two dust components exist simultaneously, were considered in [9] for the case of a Lemaitre-Tolman geometry. To model voids a spherically symmetric dust with different energy density values has been matched by an intermediate region. By introducing two dust components the boundaries no longer need to be comoving. The two-component dust model was generalized in [10] by introducing an additional homothetic vector. The obtained metrics are regular, except for a big bang singularity and in the limit of large times they can be given by an asymptotic expansion of the metric coefficients. Using this family of solutions, voids whose edge is not comoving in comparison with its surroundings can be modeled in geometric terms.

<sup>\*</sup>harko@hkucc.hku.hk<sup>†</sup>flobo@cii.fc.ul.pt

In this work, we consider the possibility that dark matter may be modeled as a mixture of two noninteracting perfect fluids, with different four-velocities. This configuration is formally equivalent to a single anisotropic fluid [11,12]. We obtain the basic equations describing the structure of this system, namely, the mass continuity and the hydrostatic equilibrium equations, as well as the equation giving the tangential velocity of test particles in stable circular orbits. For the dark matter we adopt a nonrelativistic kinetic description, with the fluid pressure proportional to the density of the fluid and to the velocity dispersion of the constituent particles. By assuming that the two fluids have the same density, the dark matter profile and the tangential velocity are obtained by numerically integrating the field equations. Some cosmological implications of the model are also discussed, and it is shown that in the large-time limit the initially anisotropic homogeneous two-fluid dark matter distribution ends in an isotropic phase.

## II. DARK MATTER AS A MIXTURE OF TWO PERFECT FLUIDS

We start our study of dark matter by assuming that it consists of a mixture of two perfect fluids, with energy densities and pressures  $\rho_1, P_1$  and  $\rho_2, P_2$ , respectively, and with four-velocities  $U^\mu$  and  $W^\mu$ , respectively. The fluid is described by the total energy-momentum tensor  $T^{\mu\nu}$ , given by

$$T^{\mu\nu} = (\rho_1 + P_1)U^\mu U^\nu - P_1 g^{\mu\nu} + (\rho_2 + P_2)W^\mu W^\nu - P_2 g^{\mu\nu}. \quad (1)$$

The four-velocities are normalized according to  $U^\mu U_\mu = 1$  and  $W^\mu W_\mu = 1$ , respectively. In the present paper we use the natural system of units with  $c = G = 1$ . The study of the physical systems described by an energy-momentum tensor having the form given by Eq. (1) can be

significantly simplified if we cast it into the standard form of perfect anisotropic fluids. This can be done by means of the transformations [11,12]

$$U^\mu \rightarrow U^{*\mu} = U^\mu \cos\alpha + \sqrt{\frac{\rho_2 + P_2}{\rho_1 + P_1}} W^\mu \sin\alpha, \quad (2)$$

$$W^\mu \rightarrow W^{*\mu} = W^\mu \cos\alpha - \sqrt{\frac{\rho_1 + P_1}{\rho_2 + P_2}} U^\mu \sin\alpha, \quad (3)$$

representing a “rotation” of the velocity four-vectors in the  $(U^\mu, W^\mu)$  velocity space, which leave the quadratic form  $(\rho_1 + P_1)U^\mu U^\nu + (\rho_2 + P_2)W^\mu W^\nu$  invariant. Thus,  $T^{\mu\nu}(U, W) = T^{\mu\nu}(U^*, W^*)$ . Next, we choose  $U^{*\mu}$  and  $W^{*\mu}$  such that one becomes timelike, while the other is spacelike. Therefore the two vectors satisfy the condition  $U^{*\mu} W^{*\mu} = 0$ . With the use of Eqs. (2) and (3) and  $U^{*\mu} W^{*\mu} = 0$  we obtain the rotation angle as

$$\tan 2\alpha = 2 \frac{\sqrt{(\rho_1 + P_1)(\rho_2 + P_2)}}{\rho_1 + P_1 - (\rho_2 + P_2)} U^\mu W_\mu. \quad (4)$$

By defining the quantities  $V^\mu = U^{*\mu}/\sqrt{U^{*\alpha}U^*_{\alpha}}$ ,  $\chi^\mu = W^{*\mu}/\sqrt{-W^{*\alpha}W^*_{\alpha}}$ ,  $\varepsilon = T^{\mu\nu}V_\mu V_\nu = (\rho_1 + P_1)U^{*\alpha}U^*_{\alpha} - (P_1 + P_2)$ ,  $\Psi = T^{\mu\nu}\chi_\mu \chi_\nu = (P_1 + P_2) - (\rho_2 + P_2)W^{*\alpha}W^*_{\alpha}$ , and  $\Pi = P_1 + P_2$ , respectively, the energy-momentum tensor of the two noninteracting perfect fluids can be written as

$$T^{\mu\nu} = (\varepsilon + \Pi)V^\mu V^\nu - \Pi g^{\mu\nu} + (\Psi - \Pi)\chi^\mu \chi^\nu, \quad (5)$$

where  $V^\mu V_\mu = 1 = -\chi^\mu \chi_\mu$  and  $\chi^\mu V_\mu = 0$  [11,12]. The energy-momentum tensor given by Eq. (5) is the standard form for anisotropic fluids [12].

The energy density  $\varepsilon$  and the radial pressure  $\Psi$  are given by

$$\varepsilon = \frac{1}{2}(\rho_1 + \rho_2 - P_1 - P_2) + \frac{1}{2}\sqrt{(\rho_1 + P_1 + \rho_2 + P_2)^2 + 4(\rho_1 + P_1)(\rho_2 + P_2)[(U^\mu W_\mu)^2 - 1]}, \quad (6)$$

$$\Psi = -\frac{1}{2}(\rho_1 + \rho_2 - P_1 - P_2) + \frac{1}{2}\sqrt{(\rho_1 + P_1 - \rho_2 - P_2)^2 + 4(\rho_1 + P_1)(\rho_2 + P_2)(U^\mu W_\mu)^2}, \quad (7)$$

respectively [11,12].

In comoving spherical coordinates  $x^0 = t$ ,  $x^1 = r$ ,  $x^2 = \vartheta$ , and  $x^3 = \phi$  we may choose  $V^1 = V^2 = V^3 = 0$ ,  $V^0 V_0 = 1$ , and  $\chi^0 = \chi^2 = \chi^3 = 0$ ,  $\chi^1 \chi_1 = -1$  [11,12]. Therefore the components of the energy-momentum of two noninteracting perfect fluids take the form:  $T^0_0 = \varepsilon$ ,  $T^1_1 = -\Psi$ , and  $T^2_2 = T^3_3 = -\Pi$ , where  $\varepsilon$  is the total energy density of the mixture of fluids,  $\Psi = P_r$  is the pressure along the radial direction, while  $\Pi = P_\perp$  is the tangential pressure on the  $r = \text{constant}$  surface.

## III. ANISOTROPIC FLUIDS IN SPHERICALLY SYMMETRIC STATIC SPACETIMES

In the following we restrict our study of the two-component dark matter to the static and spherically symmetric case, with the metric represented as

$$ds^2 = e^{\nu(r)} dt^2 - e^{\lambda(r)} dr^2 - r^2(d\vartheta^2 + \sin^2\vartheta d\phi^2). \quad (8)$$

For the metric given by Eq. (8), the Einstein gravitational field equations, describing the mixture of two fluids, can be easily integrated to give

$$e^{-\lambda} = 1 - \frac{2M(r)}{r}, \quad \text{with} \quad M(r) = 4\pi \int \varepsilon r^2 dr, \quad (9)$$

$$\frac{d\nu}{dr} = -2 \frac{\Psi'}{\varepsilon + \Psi} + \frac{4}{r} \frac{\Pi - \Psi}{\varepsilon + \Psi}, \quad (10)$$

$$\frac{dM}{dr} = 4\pi \varepsilon r^2, \quad (11)$$

$$\frac{d\Psi}{dr} = -\frac{(\varepsilon + \Psi)(4\pi\Psi r^3 + M)}{r^2(1 - 2M/r)} + \frac{2}{r}(\Pi - \Psi). \quad (12)$$

Equation (10) is the consequence of the conservation of the energy-momentum tensor,  $T_{\nu;\mu}^\mu = 0$ , while Eqs. (11) and (12) are the mass continuity, and the hydrostatic equilibrium (TOV) equation.

The galactic rotation curves provide the most direct method of analyzing the gravitational field inside a spiral galaxy. The velocity of the cloud, as measured by an inertial observer far from the source, is given by

$$v^2 = e^{-\nu} \left[ e^\lambda \left( \frac{dr}{dt} \right)^2 + r^2 \left( \frac{d\Omega}{dt} \right)^2 \right]. \quad (13)$$

For a stable circular orbit  $dr/dt = 0$ , and the tangential velocity of the test particles can be expressed as  $v_{ig}^2 = e^{-\nu} r^2 (d\Omega/dt)^2$ , where  $d\Omega^2 = d\vartheta^2 + \sin^2 \vartheta d\phi^2$ . In terms of the conserved angular momentum  $l = r^2 \dot{\phi}$  and energy  $E = e^{\nu(r)} \dot{t}$ , the tangential velocity is given, for  $\theta = \pi/2$ , by  $v_{ig}^2 = e^\nu l^2 / r^2 E^2$ . By taking into account the explicit expressions for  $l$  and  $E$  we obtain for the tangential velocity of a test particle in a stable circular orbit the expression  $v_{ig}^2 = r\nu'/2$ .

The tangential velocity profile can be obtained as a function of the total density and of the mass of the dark matter from Eq. (10) as

$$v_{ig}^2(r) = \frac{4\pi\Psi r^3 + M}{r(1 - 2M/r)}. \quad (14)$$

Equation (12), can be written in the following equivalent form:

$$\frac{d\Psi}{dr} = -\frac{(\varepsilon + \Psi)}{r} v_{ig}^2(r) + \frac{2}{r}(\Pi - \Psi). \quad (15)$$

#### IV. DARK MATTER AS A MIXTURE OF TWO FLUIDS

The kinetic energy-momentum tensor  $T_\nu^\mu$  associated to the frozen distribution of dark matter is given by  $T_\nu^\mu = (g/h^3) \int d^3p f(p) p^\mu p_\nu / p^0$ , where  $f(p)$  is the dark matter particle distribution function,  $p^\mu$  is the four-momentum,  $\vec{p}$  is the three-momentum, with absolute value  $p$ , and  $g$  is the number of helicity states, respectively. The energy density  $\varepsilon$  of the system is defined as

$$\varepsilon = \frac{g}{3h^3} \int E f(p) d^3p, \quad (16)$$

while the pressure of a system with an isotropic distribution of momenta is given by

$$P = \frac{g}{3h^3} \int p v f(p) d^3p = \frac{g}{3h^3} \int \frac{p^2}{E} f(p) d^3p, \quad (17)$$

where the velocity  $v$  is related to the momentum by  $v = p/E$  [13]. In the nonrelativistic regime, when  $E \approx m$  and  $p \approx mv$ , the density  $\rho$  of the dark matter is given by  $\rho = mn$ , where  $n$  is the particle number density, while its pressure  $P$  can be obtained as [13]

$$P = \frac{g}{3h^3} \int \frac{p^2 c^2}{E} f(p) d^3p \approx 4\pi \frac{g}{3h^3} \int \frac{p^4}{m} dp, \quad (18)$$

giving

$$P = \sigma^2 \rho, \quad (19)$$

where  $\sigma^2 = \langle \vec{v}^2 \rangle / 3$ , and  $\langle \vec{v}^2 \rangle$  is the average squared velocity of the particle.  $\sigma$  is the one-dimensional velocity dispersion. In the nonrelativistic approximation given by Eqs. (18) and (19), the velocity dispersion  $\sigma$  is a constant only for the case of the nondegenerate ideal Maxwell-Boltzmann gas, whose Newtonian analogue is the isothermal sphere. In the following we will consider only the specific case of the ideal gases in the nonrelativistic regime, and we will explicitly assume that the dispersion velocity is a constant.

We model the galactic dark matter as a mixture of two fluids, with energy densities  $\rho_i$ ,  $i = 1, 2$ , constant velocity dispersions  $\sigma_i$ ,  $i = 1, 2$ , pressures  $P_i$ ,  $i = 1, 2$ , and with four-velocities  $U^\mu$  and  $W^\mu$ , respectively. Since the two fluids have different four-velocities, we may write  $U^\mu W_\mu = 1 + b/2$ , where generally  $b$  is an arbitrary function of the radial coordinate. The functional form of  $b$  can be obtained from Eq. (4), which gives

$$U^\mu W_\mu = \frac{1}{2} \tan 2\alpha \left[ \sqrt{\frac{\rho_1 + p_1}{\rho_2 + p_2}} - \sqrt{\frac{\rho_2 + p_2}{\rho_1 + p_1}} \right], \quad (20)$$

while  $b$  can be obtained as

$$b = \tan 2\alpha \left[ \sqrt{\frac{\rho_1 + p_1}{\rho_2 + p_2}} - \sqrt{\frac{\rho_2 + p_2}{\rho_1 + p_1}} \right] - 2. \quad (21)$$

In the following we choose our coordinates such that we are comoving with the fluid with velocity  $U^\mu$ . Thus  $U^0 = e^{-\nu/2}$  and  $U^1 = U^2 = U^3 = 0$ . We also assume  $W^2 = W^3 = 0$ . From the conditions  $U_\mu U^\mu = 1$  and  $W_\mu W^\mu = 1$  we obtain  $W^0 = (1 + b/4)e^{-\nu/2}$  and  $W^1 = \sqrt{(b/2)(1 + b/8)}e^{-\lambda/2}$ . By assuming that the physical parameters of the two fluids satisfy the condition

$$2b \left( 1 + \frac{b}{8} \right) \frac{(\rho_1 + p_1)(\rho_2 + p_2)}{(\rho_1 + \rho_2 + p_1 + p_2)^2} \ll 1, \quad (22)$$

the energy density, and the radial and tangential pressures can be obtained as

$$\varepsilon = \rho_1 + \rho_2 + \frac{b}{2} \left( 1 + \frac{b}{8} \right) \frac{(\rho_1 + P_1)(\rho_2 + P_2)}{(\rho_1 + \rho_2 + P_1 + P_2)}, \quad (23)$$

$$\Psi = P_1 + P_2 + \frac{b}{2} \left( 1 + \frac{b}{8} \right) \frac{(\rho_1 + P_1)(\rho_2 + P_2)}{(\rho_1 + \rho_2 + P_1 + P_2)}, \quad (24)$$

$$\Pi = P_1 + P_2. \quad (25)$$

In the following we assume that the masses  $m_1$  and  $m_2$  and the number densities  $n_1$  and  $n_2$  of the dark matter particles satisfy the scaling relation  $m_1/m_2 = n_2/n_1$ , giving  $\rho_1 = \rho_2 = \rho$  for the densities of the two components. Taking into account that  $P_1 = \sigma_1^2 \rho_1$ ,  $P_2 = \sigma_2^2 \rho_2$ , the expression of  $b$  can be written as

$$b = \tan 2\alpha \left( \frac{\sigma_1^2 - \sigma_2^2}{\sqrt{1 + \sigma_1^2} \sqrt{1 + \sigma_2^2}} \right) - 2. \quad (26)$$

Hence for this particular case it follows that  $b = \text{constant}$ . Moreover, since the rotation angle is arbitrary, for  $\sigma_1 > \sigma_2$ , we can always find an angle  $\alpha$  so that the condition  $b > 0$  is also satisfied. In the following we will restrict our analysis to the range of positive values of  $b$ .

Therefore the energy density and the pressures of the two-fluid dark matter mixture can be written as  $\varepsilon = \varepsilon_0 \rho$ ,  $\Psi = \Psi_0 \rho$ ,  $\Pi = \Pi_0 \rho$ , where  $\varepsilon_0 = [2 + (b/2)(1 + b/8)]$ ,  $\Psi_0 = [\sigma_1^2 + \sigma_2^2 + (b/2)(1 + b/8)]$ , and  $\Pi_0 = \sigma_1^2 + \sigma_2^2$ , respectively. By introducing a set of dimensionless variables  $(\theta, \xi, m)$ , defined as  $\varepsilon = \varepsilon_c \theta$ ,  $r = \xi / \sqrt{4\pi\varepsilon_c}$ ,  $M = m / \sqrt{4\pi\varepsilon_c}$ , where  $\varepsilon_c$  is the central density of the dark matter halo, the continuity equation, the TOV equation and the tangential velocity can be written in a dimensionless form as

$$\frac{dm}{d\xi} = \xi^2 \theta, \quad (27)$$

$$\frac{d\theta}{d\xi} = - \frac{(1 + \gamma_\varepsilon) \theta (\xi^3 \theta / \gamma_\varepsilon + m)}{\xi^2 (1 - 2m/\xi)} + 2(\gamma_\Pi - 1) \frac{\theta}{\xi}, \quad (28)$$

$$v_{tg}^2 = \frac{\xi^3 \theta / \gamma_\varepsilon + m}{\xi (1 - 2m/\xi)}, \quad (29)$$

where  $\gamma_\varepsilon = \varepsilon_0 / \Psi_0$  and  $\gamma_\Pi = \Pi_0 / \Psi_0$ , respectively. Eqs. (27) and (28) must be integrated with the boundary conditions  $\theta(0) = 1$  and  $m(0) = 1$ . The variation of the two-fluid dark matter density profile and of the tangential velocity of test bodies is represented, for fixed values of  $\sigma_1$  and  $\sigma_2$ , in Figs. 1 and 2.

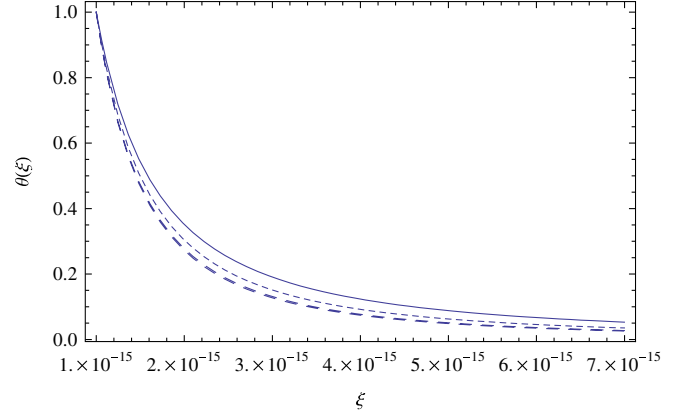


FIG. 1 (color online). Dimensionless density profile of a two-component dark matter fluid, with  $\sigma_1 = 950$  km/s and  $\sigma_2 = 750$  km/s, for different values of  $b$ :  $b = 0.0001$  (solid curve),  $b = 0.0002$  (dotted curve),  $b = 0.0004$  (dashed curve), and  $b = 0.0005$  (long-dashed curve), respectively.

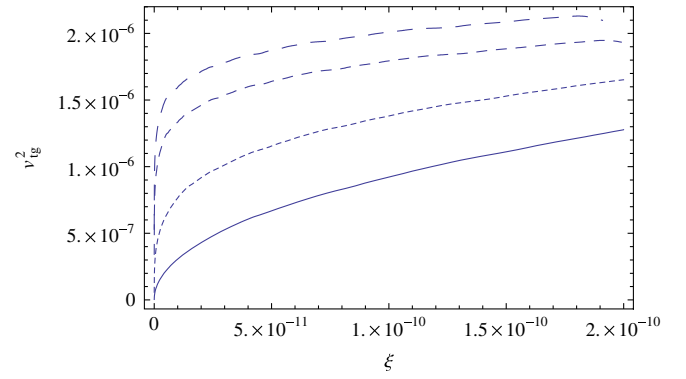


FIG. 2 (color online). Tangential velocities of test bodies in a two-component dark matter fluid, with  $\sigma_1 = 950$  km/s and  $\sigma_2 = 750$  km/s, for different values of  $b$ :  $b = 0.0001$  (solid curve),  $b = 0.0002$  (dotted curve),  $b = 0.0004$  (dashed curve), and  $b = 0.0005$  (long-dashed curve), respectively.

## V. COSMOLOGICAL IMPLICATIONS

Since the two-fluid dark matter mixture represents an anisotropic fluid, on a cosmological scale the corresponding geometry is also anisotropic. In a homogenous universe a two-fluid mixture can be described by a flat Bianchi type I geometry, with the line element given by

$$ds^2 = dt^2 - a_1^2(t) dx^2 - a_2^2(t) dy^2 - a_3^2(t) dz^2, \quad (30)$$

where  $a_i(t)$ ,  $i = 1, 2, 3$  are the directional scale factors. In this geometry the gravitational field equations take the form

$$3\dot{H} + H_1^2 + H_2^2 + H_3^2 = -\frac{1}{2}(\varepsilon + \Psi + 2\Pi), \quad (31)$$

$$\frac{1}{V} \frac{d}{dt} (V H_1) = \frac{1}{2}(\varepsilon - \Psi), \quad (32)$$



$$\frac{1}{V} \frac{d}{dt}(VH_2) = \frac{1}{V} \frac{d}{dt}(VH_3) = \frac{1}{2}(\varepsilon - \Pi), \quad (33)$$

where we have denoted  $V = a_1 a_2 a_3$ ,  $H_i = \dot{a}_i/a_i$ ,  $i = 1, 2, 3$  and  $H = \frac{1}{3} \sum_{i=1}^3 H_i = \dot{V}/3V$ , respectively. On a cosmological scale the pressures of the two-fluid dark matter obey the conditions  $\Psi, \Pi \ll \varepsilon$ , and hence in the gravitational field equations we can neglect  $\Psi$  and  $\Pi$  with respect to  $\varepsilon$ . By adding Eqs. (32) we obtain

$$\dot{H} + 3H^2 = \frac{1}{2}\varepsilon, \quad (34)$$

leading to  $H_i = H + K_i/V$ ,  $i = 1, 2, 3$ , where  $K_i$ ,  $i = 1, 2, 3$  are constants of integration satisfying the consistency condition  $\sum_{i=1}^3 K_i = 0$ . We also denote  $K^2 = \sum_{i=1}^3 K_i^2$ . From the conservation of the energy density of the dark matter,  $\dot{\varepsilon} + 3H\varepsilon = 0$ , we obtain  $\varepsilon = \varepsilon_C/V$ , where  $\varepsilon_C$  is an arbitrary integration constant. Equation (34) takes the form  $\dot{V} = (3/2)\varepsilon_C$ , with the general solution given by  $V(t) = (3/4)\varepsilon_C t^2 + C_1 t + C_2$ , with  $C_i$ ,  $i = 1, 2$  arbitrary integration constant. The large-time behavior of an anisotropic cosmological model can be obtained from the study of the anisotropy parameter  $A$ , defined by

$$A = \frac{1}{3} \sum_{i=1}^3 \left( \frac{H_i - H}{H} \right)^2 = \frac{3K^2}{V^2}. \quad (35)$$

If  $A = 0$  the cosmological model is isotropic. For the two-fluid dark matter model  $A = 3K^2/[(3/2)\varepsilon_C t + C_1]^2$ , and therefore in the large-time limit  $A \rightarrow 0$ . Thus, on a cosmological scale, and in the large-time limit, the initially anisotropic dark matter fluid mixture behaves like an isotropic fluid.

## VI. DISCUSSIONS AND FINAL REMARKS

In the present paper we have considered the theoretical possibility that dark matter may be modeled as a mixture of two noninteracting perfect fluids, with different four-velocities. This model is formally equivalent to a single anisotropic fluid. By considering a purely kinetic description of dark matter the basic equations describing the dark matter mixture have been obtained, and integrated numerically. In the particular case in which  $\Psi, \Pi \ll \varepsilon$ , we can

approximate  $\Psi = \Pi \approx 0$ , corresponding to two noninteracting pressureless fluids. For this case the theoretical model reduces to a single anisotropic fluid with vanishing tangential pressure. In the general case of arbitrary fluids there is a general  $r$ -dependent functional relationship between the energy density and the radial and tangential pressures, which differs radically from the simple barotropic equation of state. On the other hand, it was shown in [8] that half of the 14 galaxies considered in the study are well fit by the polytropic halo model, despite its serious physical simplifications. The necessity of considering non-standard dark matter models with pressure is justified by our uncertainties in the knowledge about the nature of the dark matter particle, as well as by the fact that this model gives a much better description of the observational results, as compared to the pressureless case. As a consistency check of our model we have considered the cosmological behavior of a mixture of two fluids, and we have showed that such a mixtures always isotropizes in the large-time limit.

The approach used in the present paper is different from the method used in [9] for the study of dust shells. First of all, the two fluids are described as a single effective anisotropic fluid. While in [9] a noncomoving frame of reference is adopted, due to our transformation of the energy-momentum tensor, we can write down and study the field equations in a comoving frame. Moreover, the inclusion of the pressure terms is straightforward. We have also to point out the differences in geometry: while in [9] the time-dependent Lemaitre-Tolman model is used, and additional symmetries are imposed in order to reduce the number of variables, in the present paper we are considering the properties of dark matter in a spherically symmetric static geometry. The present model can be easily generalized for other equations of state of dark matter, or for cases involving the presence of condensates or scalar fields. These possibilities, and their impact on dark matter distribution, as well as some cosmological implications, will be considered in some future studies.

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- [1] P.J.E. Peebles and B. Ratra, *Rev. Mod. Phys.* **75**, 559 (2003); T. Padmanabhan, *Phys. Rep.* **380**, 235 (2003).  
[2] J.M. Overduin and P.S. Wesson, *Phys. Rep.* **402**, 267 (2004).

- [3] U. Nucamendi, M. Salgado, and D. Sudarsky, *Phys. Rev. Lett.* **84**, 3037 (2000); T. Matos and F.S. Guzman, *Classical Quantum Gravity* **18**, 5055 (2001); E.W. Mielke and F.E. Schunk, *Phys. Rev. D* **66**, 023503 (2002); B. Fuchs and E.W. Mielke, *Mon. Not. R.*

- Astron. Soc.* **350**, 707 (2004); X. Hernández, T. Matos, R. A. Sussman, and Y. Verbin, *Phys. Rev. D* **70**, 043537 (2004); D. Giannios, *Phys. Rev. D* **71**, 103511 (2005); A. Bernal and F. S. Guzman, *Phys. Rev. D* **74**, 063504 (2006); C. G. Boehmer and T. Harko, *J. Cosmol. Astropart. Phys.* **06** (2007) 025; T. Harko, *Mon. Not. R. Astron. Soc.* **413**, 3095 (2011); P.-H. Chavanis, [arXiv:1103.2050](#); P.-H. Chavanis and L. Delfini, [arXiv:1103.2054](#); P.-H. Chavanis, [arXiv:1103.2698](#); P.-H. Chavanis, [arXiv:1103.3219](#); F. Briscese, *Phys. Lett. B* **696**, 315 (2011); T. Harko, *J. Cosmol. Astropart. Phys.* **05** (2011) 022.
- [4] S. Capozziello, V. F. Cardone, and A. Troisi, *J. Cosmol. Astropart. Phys.* **08** (2006) 001; *Mon. Not. R. Astron. Soc.* **375**, 1423 (2007); A. Borowiec, W. Godlowski, and M. Szydlowski, *Int. J. Geom. Methods Mod. Phys.* **4**, 183 (2007); C. F. Martins and P. Salucci, *Mon. Not. R. Astron. Soc.* **381**, 1103 (2007); C. G. Boehmer, T. Harko, and F. S. N. Lobo, *Astropart. Phys.* **29**, 386 (2008); *J. Cosmol. Astropart. Phys.* **03** (2008) 024; O. Bertolami, C. G. Boehmer, T. Harko, and F. S. N. Lobo, *Phys. Rev. D* **75**, 104016 (2007); F. S. N. Lobo, [arXiv:0807.1640](#).
- [5] J. F. Navarro, C. S. Frenk, and S. D. M. White, *Mon. Not. R. Astron. Soc.* **275**, 720 (1995).
- [6] A. Burkert, *Astrophys. J.* **447**, L25 (1995).
- [7] S. Bharadwaj and S. Kar, *Phys. Rev. D* **68**, 023516 (2003); K.-Y. Su and P. Chen, *Phys. Rev. D* **79**, 128301 (2009).
- [8] C. J. Saxton and I. Ferreras, *Mon. Not. R. Astron. Soc.* **405**, 77 (2010).
- [9] G. Haager, *Classical Quantum Gravity* **14**, 2219 (1997).
- [10] G. Haager, *Classical Quantum Gravity* **15**, 3669 (1998).
- [11] P. S. Letelier, *Phys. Rev. D* **22**, 807 (1980); P. S. Letelier and P. S. C. Alencar, *Phys. Rev. D* **34**, 343 (1986).
- [12] L. Herrera and N. O. Santos, *Phys. Rep.* **286**, 53 (1997).
- [13] C. J. Hogan and J. J. Dalcanton, *Phys. Rev. D* **62**, 063511 (2000); J. Madsen, *Phys. Rev. D* **64**, 027301 (2001).